

On
Volume-Surface Reaction-Diffusion systems

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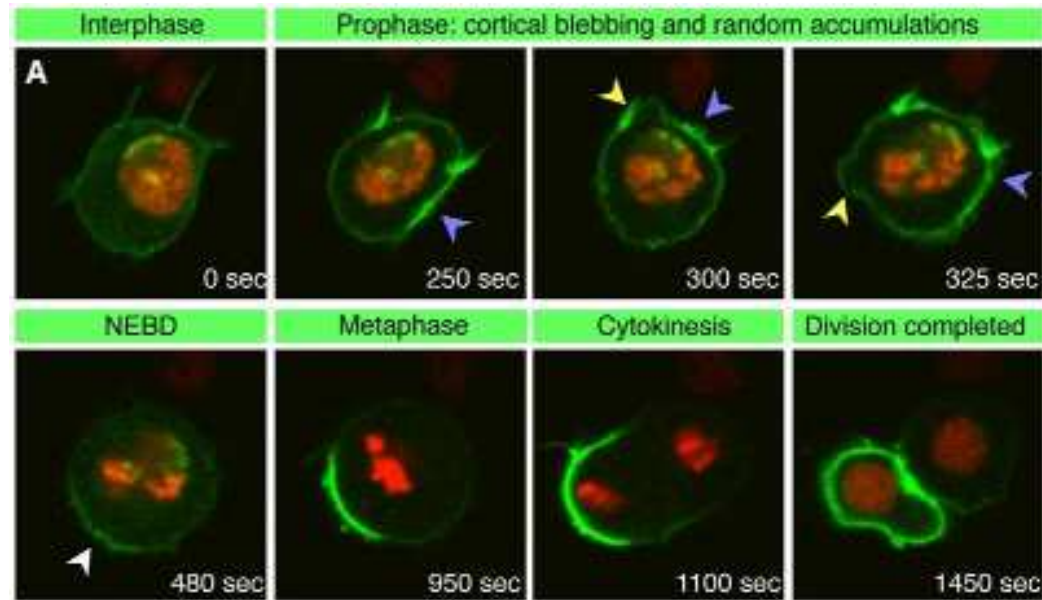
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joint works with L. Desvillettes, H. Egger, J.-F. Pietschmann,
E. Latos, B.Q. Tang

Complex-Balanced Volume-Surface RD Network

Protein-localisation before asymmetric stem-cell division



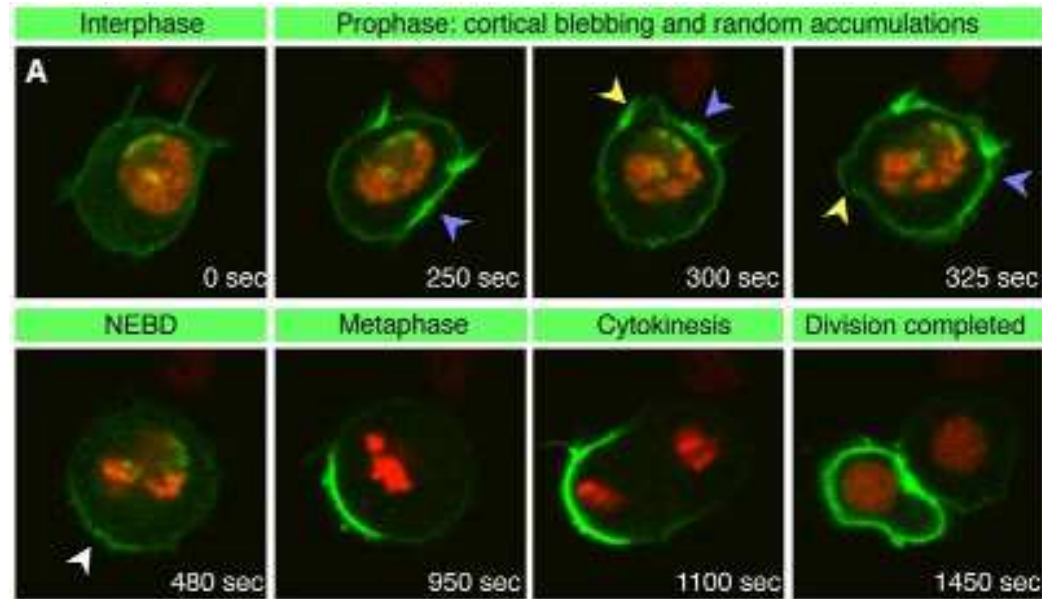
Asymmetric stem-cell division:

Cell-diversity by localisation of cell-fate determinants into one side of the cell cortex and into one of two daughter cells.^a

^aGFP-Pon in SOP precursor cells in living *Drosophila* larvae [Meyer, Emery, Berdnik, Wirtz-Peitz, Knoblich, *Current Biology*, 2005]

Complex-Balanced Volume-Surface RD Network

Protein-localisation before asymmetric stem-cell division



Mathematical model:

- “high” concentrations, insignificant stochastic effects
- system of (reversible) reaction-diffusion equations
- volume(cytoplasm)-surface(membran) dynamics

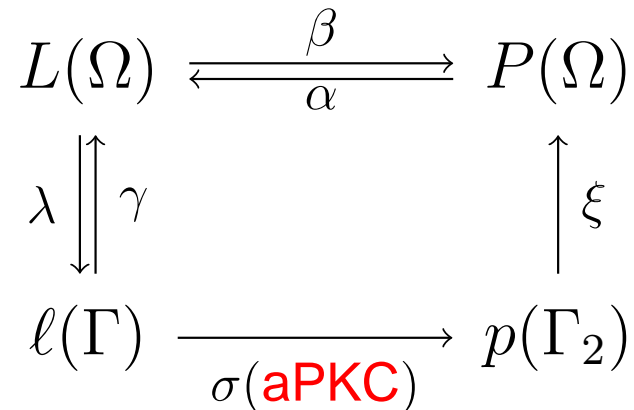
A Volume-Surface Reaction-Diffusion Model

Model Assumptions and Quantities



Key protein: Lgl in cytoplasm (Ω) and cell cortex ($\Gamma = \partial\Omega$).

Key kinase: aPKC phosphorylates Lgl on a part Γ_2 of cortex.



$L(t, x)$ cytoplasmic Lgl \leftrightarrow $l(t, x)$ cortical Lgl \rightarrow activation of aPKC

$\rightarrow p(t, x)$ cortical p-Lgl $\rightarrow P(t, x)$ cytoplasmic p-Lgl $\leftrightarrow L(t, x)$

Complex-balanced reaction-diffusion network

Bio: qualitative interplay reaction/surface/volume diffusion

A Volume-Surface Reaction-Diffusion Model

Detailed versus Complex Balance Equilibria



A detailed balance equilibrium **balances the forward and backward reactions between** all species/complexes.

A complex balance equilibrium **balances the total inflow and total outflow** from and into all species/complexes.

A Volume-Surface Reaction-Diffusion Model

A prototypical model I



Volume equations with diffusion coefficients $d_L, d_P > 0$

$$(V) \quad \begin{cases} L_t - d_L \Delta L = \alpha P - \beta L, & x \in \Omega, t > 0, \\ P_t - d_P \Delta P = -\alpha P + \beta L, & x \in \Omega, t > 0, \\ L(0, x) = L_0(x), P(0, x) = P_0(x), & x \in \Omega \end{cases}$$

Boundary conditions on $\partial\Omega = \Gamma = \Gamma_1 \cup \Gamma_2$ and $\Gamma_1 \cap \Gamma_2 = \emptyset$

$$(BC) \quad \begin{cases} d_L \frac{\partial L}{\partial \nu} = \gamma l - \lambda L, & x \in \Gamma, t > 0, \\ d_P \frac{\partial P}{\partial \nu} = 0, & x \in \Gamma_1, t > 0, \\ d_P \frac{\partial P}{\partial \nu} = \xi p, & x \in \Gamma_2, t > 0, \end{cases}$$

Reaction rates $\alpha, \beta, \gamma, \lambda, \sigma, \xi$ are positive constants

A Volume-Surface Reaction-Diffusion Model

A prototypical model II



Boundary dynamics

$$(\text{BD}) \quad \left\{ \begin{array}{ll} l_t - d_l \Delta_{\Gamma} l = \lambda L - \gamma l - \sigma \chi_{\Gamma_2} l, & x \in \Gamma, t > 0 \\ p_t - d_p \Delta_{\Gamma_2} p = \sigma l - \xi p, & x \in \Gamma_2, t > 0, \\ d_p \frac{\partial p}{\partial \nu_{\Gamma_2}} = 0, & x \in \partial \Gamma_2, \\ l(0, x) = l_0(x), & x \in \Gamma, \\ p(0, x) = p_0(x), & x \in \Gamma_2, \end{array} \right.$$

Δ is the usual **Laplacian** in the domain Ω

Δ_{Γ} and Δ_{Γ_2} are **Laplace-Beltrami** operator on Γ and Γ_2

χ_{Γ_2} is the characteristic function of Γ_2

A Volume-Surface Reaction-Diffusion Model

Properties and Local well-posedness



Conservation law: total Lgl mass

$$\frac{d}{dt} \left[\int_{\Omega} (L(t, x) + P(t, x)) + \int_{\Gamma} l(t, x) + \int_{\Gamma_2} p(t, x) \right] = 0.$$

Local well-posedness:

There exists of a **unique weak/strong local solution** (L, P, l, p) on $(0, T)$, which is **non-negative** if the initial data are so.^a

^a[K.F., S. Rosenberger, B.Q. Tang, Comm. Math. Sciences 2016]

A Volume-Surface Reaction-Diffusion Model



Complex balance reaction network

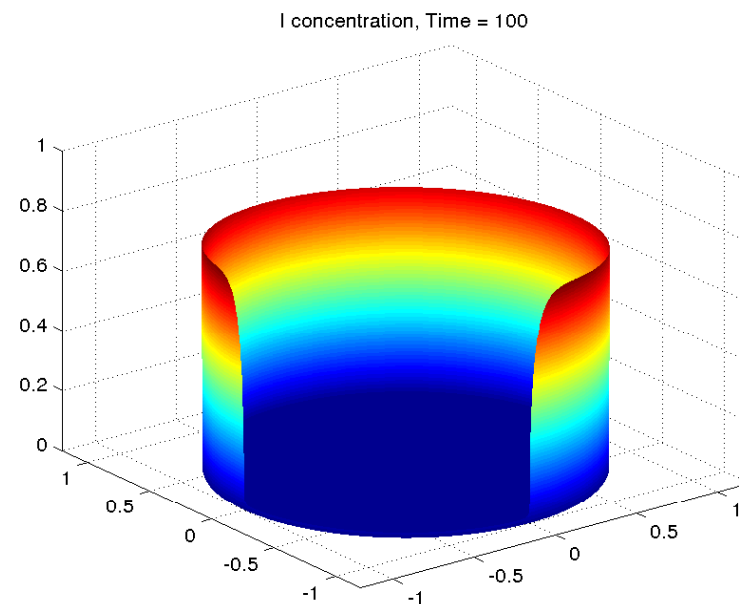
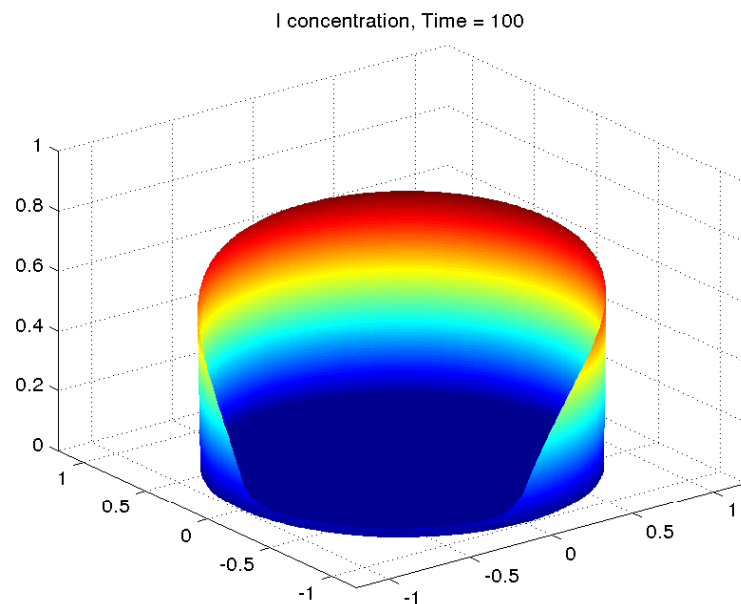


Figure 1: l -Lgl(Γ) with and without surface diffusion

Numerical analysis of VSRD models including **discrete entropy structure/estimates**: ^a

^a[Egger, F., Pietschmann, Tang, to appear in Applied Math & Computation]

A Volume-Surface Reaction-Diffusion Model

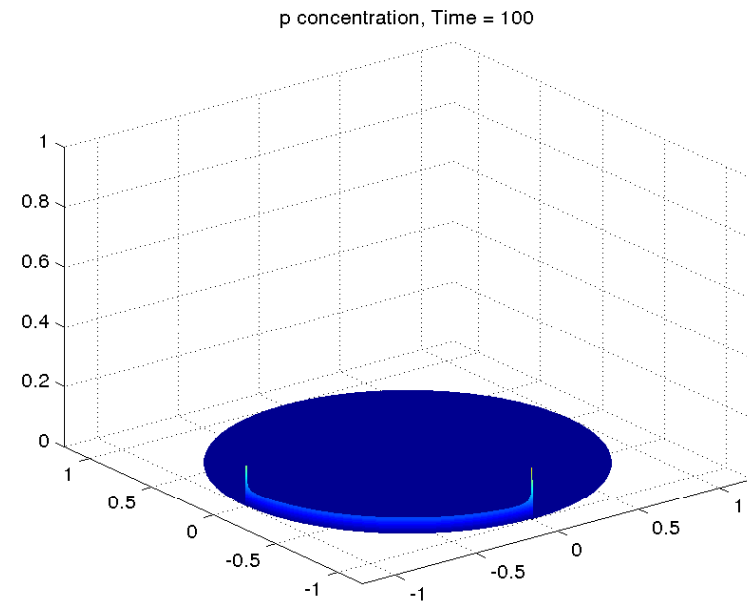
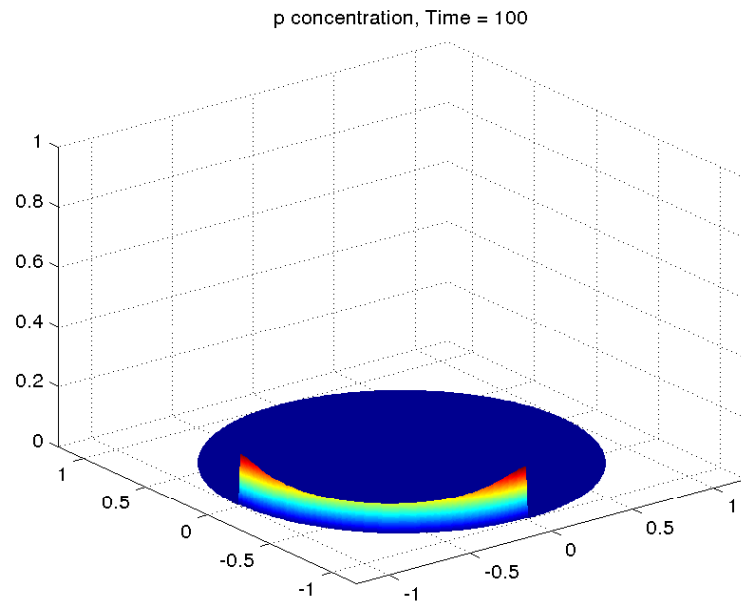


Figure 2: p -Lgl(Γ) with and without surface diffusion

Surface diffusion $O(10^{-2})$: indirect surface diffusion effect via weakly reversible reaction $O(1)$ and volume diffusion $O(10^{-2})$

A Volume-Surface Reaction-Diffusion Model

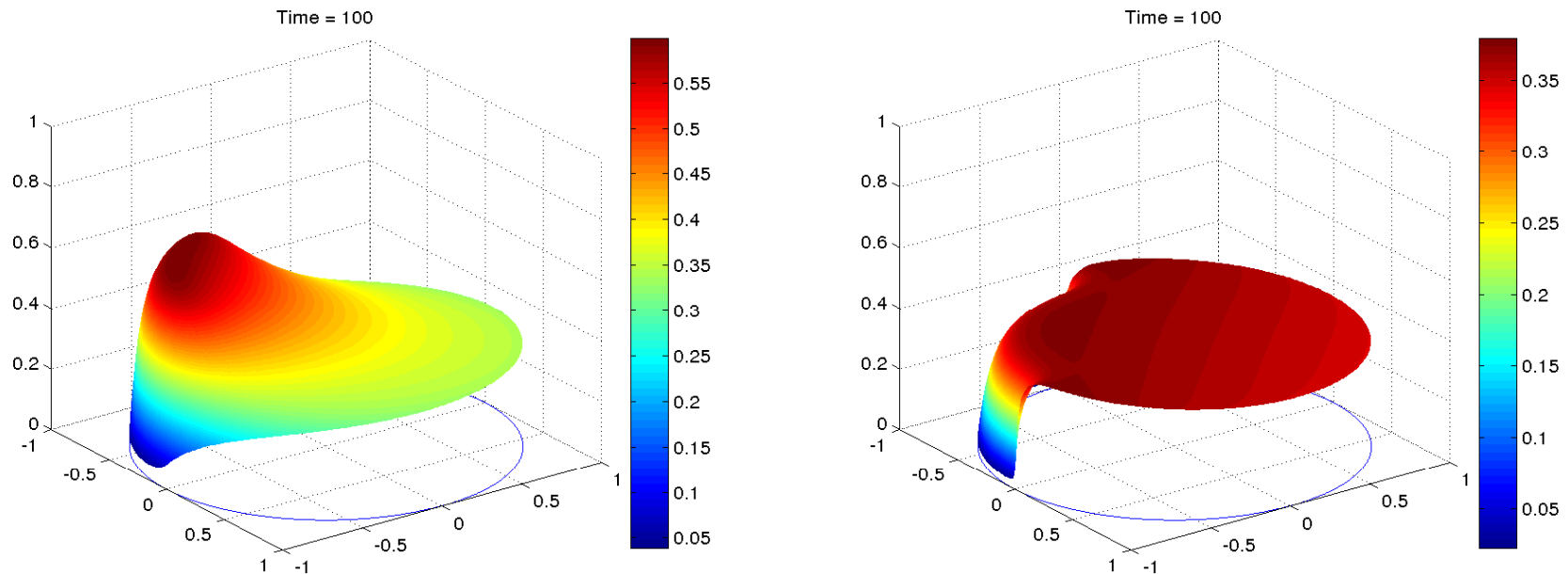


Figure 3: L -Lgl(Ω) with and without surface diffusion

Surface diffusion and weakly reversible reaction lead to stationary hump in L within Ω .

A Volume-Surface Reaction-Diffusion Model

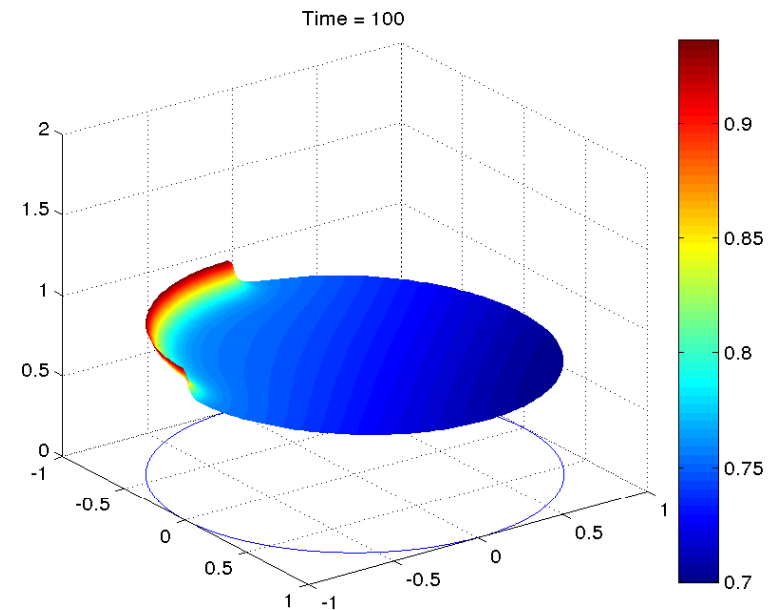
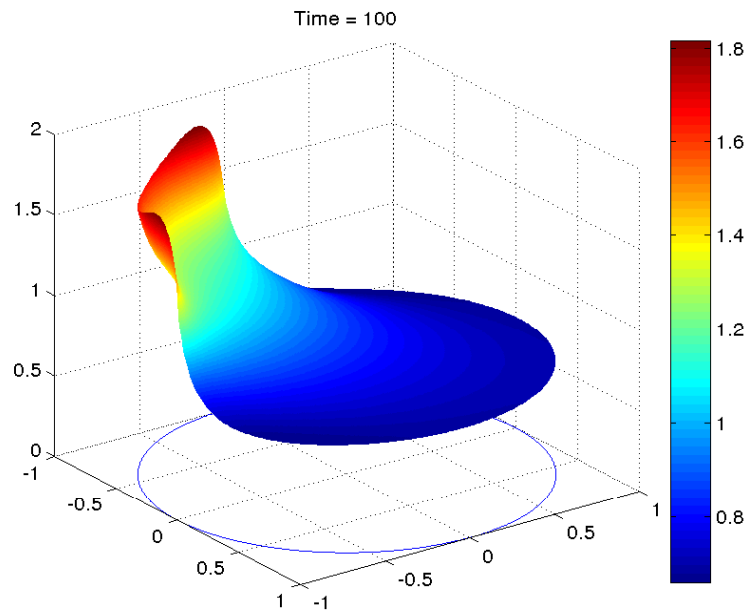


Figure 4: P -Lgl(Ω) with and without surface diffusion

Stationary hump in L as consequence of inflow from p into $P \rightarrow L$ and shape of Ω .

A Volume-Surface Reaction-Diffusion Model

Global existence and large time behaviour



Theorem: Unique global-in-time weak solution (L, P, l, p) .

Proof: L^2 -type energy estimate and Gronwall.

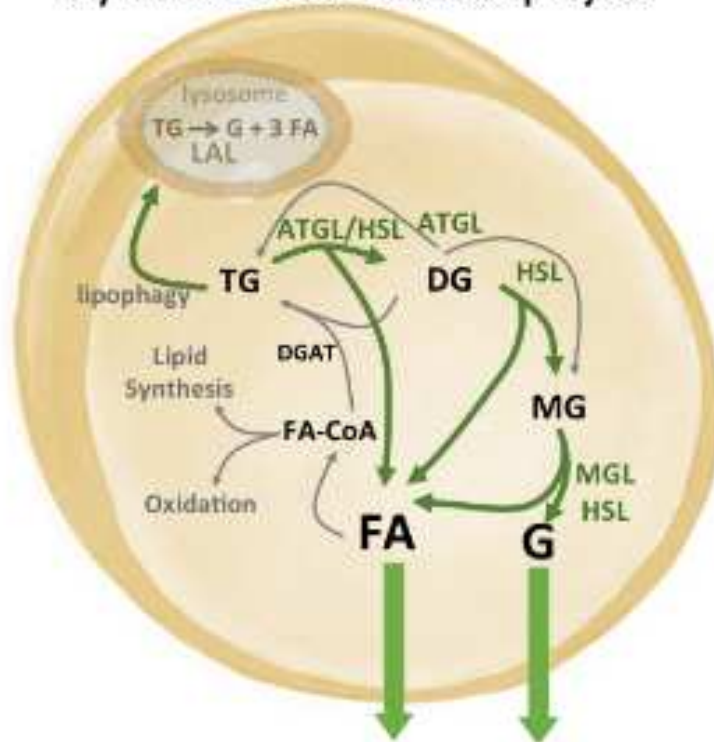
Question:

- Convergence to complex balance equilibrium for all initial data and parameter?
- L^2 -Entropy?

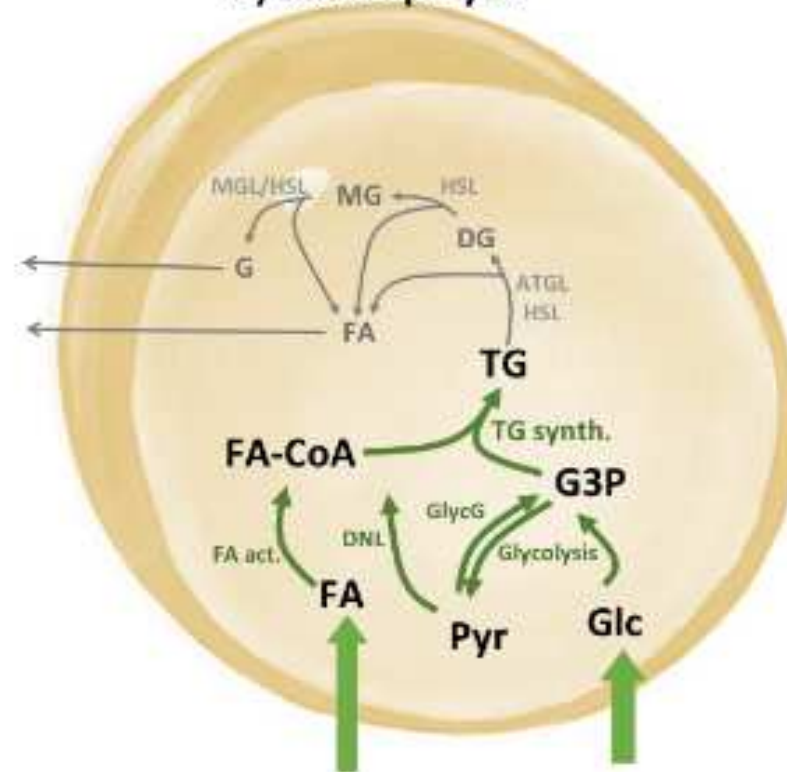
Another (Volume-Surface) RD Model

Lipolysis

A) hormone-stimulated lipolysis



B) basal lipolysis



Lipolysis: Breakdown of lipids and hydrolysis of triglycerides into glycerol and fatty acids.

Systems of Reaction-Diffusion Equations

Nonlinear Complex Balance Networks



Substances: $\mathcal{S} = \{S_1, \dots, S_N\}$,

Complexes: $\mathcal{C} = \{\mathbf{y}_1, \dots, \mathbf{y}_{|\mathcal{C}|}\}$ with $\mathbf{y}_i \in (\{0\} \cup [1, \infty))^N$,

Reactions: $\mathcal{R} = \{\mathbf{y} \rightarrow \mathbf{y}'\}$ from source \mathbf{y} into product $\mathbf{y}' \in \mathcal{C}$.

Mass action law reaction rate for $\mathbf{y}_r \rightarrow \mathbf{y}'_r$: $\mathbf{c}^{\mathbf{y}_r} = \prod_{i=1}^N c_i^{y_{r,i}}$

Reaction rate constant k_r of the reaction $\mathbf{y}_r \rightarrow \mathbf{y}'_r$.

Reaction vector: $\mathbf{R}(\mathbf{c}) = \sum_{r=1}^{|\mathcal{R}|} k_r \mathbf{c}^{\mathbf{y}_r} (\mathbf{y}'_r - \mathbf{y}_r)$

Nonlinear reaction-diffusion network

$$\frac{\partial}{\partial t} \mathbf{c} - \mathbb{D} \Delta \mathbf{c} = \mathbf{R}(\mathbf{c}) \quad \text{for } (x, t) \in \Omega \times (0, +\infty),$$

with $\mathbb{D} = \text{diag}(d_1, \dots, d_N)$.

Homogeneous Neumann BCs on Lipschitz domain Ω .

[JH72]: A **complex balanced** network has a **unique positive equilibrium**, which **balances the total outflow and inflow** for all complexes $\mathbf{y} \in \mathcal{C}$:

$$\sum_{\{r: \mathbf{y}_r = \mathbf{y}\}} k_r \mathbf{c}_\infty^{\mathbf{y}_r} = \sum_{\{s: \mathbf{y}'_s = \mathbf{y}\}} k_s \mathbf{c}_\infty^{\mathbf{y}'_s}.$$

Relative (free energy) entropy functional

$$\mathcal{E}(\mathbf{c}|\mathbf{c}_\infty) = \sum_{i=1}^N \int_{\Omega} \left(c_i \log \frac{c_i}{c_{i,\infty}} - c_i + c_{i,\infty} \right) dx$$

Explicit (**nontrivial**) entropy dissipation functional with

$$e(x, y) = x \log(x/y) - x + y$$

$$\begin{aligned} \mathcal{D}(\mathbf{c}) &= -\frac{d}{dt} \mathcal{E}(\mathbf{c}|\mathbf{c}_\infty) \\ &= \sum_{i=1}^N d_i \int_{\Omega} \frac{|\nabla c_i|^2}{c_i} dx + \sum_{r=1}^{|\mathcal{R}|} k_r \mathbf{c}_\infty^{y_r} e \left(\frac{\mathbf{c}^{y_r}}{\mathbf{c}_\infty^{y_r}}, \frac{\mathbf{c}^{y'_r}}{\mathbf{c}_\infty^{y'_r}} \right) \geq 0 \end{aligned}$$

Theorem:^a For **complex balanced RD networks without boundary equilibria**, any renormalised (Fisher [2015]) solution $\mathbf{c}(x, t)$ converges exponentially to \mathbf{c}_∞ in L^1 with a rate $\lambda/2$:

$$\sum_{i=1}^N \|c_i(t) - c_{i,\infty}\|_{L^1(\Omega)}^2 \leq C_{\text{CKP}}^{-1} \mathcal{E}(\mathbf{c}_0 | \mathbf{c}_\infty) e^{-\lambda t} \quad \text{for a.a. } t > 0,$$

where C_{CKP} is the constant in a Csiszár-Kullback-Pinsker type inequality.

Renormalised solutions satisfy all **mass/charge conservation laws** and a **weak entropy-dissipation law**, Fisher [2017]

^a[K.F. B.Q.Tang, ZAMP 2018]

The Entropy Method



Quantitative large-time behaviour

$\mathcal{E}(f)$ non-increasing **convex** entropy functional

$\mathcal{D}(f)$ entropy production, f_∞ entropy minimising equilibrium

$$\frac{d}{dt}\mathcal{E}(f) = \frac{d}{dt}\mathcal{E}(f) - \mathcal{E}(f_\infty) = -\mathcal{D}(f) \leq 0$$

provided **conservation laws**: $\mathcal{D}(f) = 0 \iff f = f_\infty$

$$\mathcal{D} \geq \Phi(\mathcal{E}(f) - \mathcal{E}(f_\infty)), \quad \Phi(0) = 0, \quad \Phi \geq 0$$

\Rightarrow **explicit convergence in entropy**, exponential if $\Phi'(0) > 0$

\Rightarrow convergence in L_1 : $\|f - f_\infty\|_1^2 \leq C(\mathcal{E}(f) - \mathcal{E}(f_\infty))$

Csiszár-Kullback-Pinsker inequalities for convex entropies

Entropy Method

Advantages:

- based on **functional inequalities** → "robust"
- avoids linearisation → "global" results
- allows for **explicit** constants

nonlinear diffusion: [T], [CJMTU], [AMTU], [DV]...

inhomogeneous kinetic equations: [DV], ...

reaction-diffusion systems: [Grö83], [Grö92], [DF06], [DF08],
[DF14], [MMH15], [FL16], [PSZ17], [DFT17], [FT17],
[HHMM18], [FT18] **no Bakry-Emery strategy**



Theorem:^a For any **complex balanced reaction networks without boundary equilibria**, there exists a constant $\lambda > 0$ and the “exponential” **entropy entropy-dissipation estimate**

$$\mathcal{D}(\mathbf{c}(t)) \geq \lambda \mathcal{E}(\mathbf{c}(t) | \mathbf{c}_\infty),$$

- Proof via convexification: [MMH15], [PSZ17] (detailed b.)
- Proof via explicit estimates using conservation laws
 $\mathbb{Q} \bar{\mathbf{c}} = \mathbf{M}$: [DFT17], [FT17], [FLT18]
Method applies also to volume-surface RD systems
- Proof via reduction to finite-dimensional inequality: [FT18]

^a[L. Desvillettes, K.F., B.Q. Tang, SIMA 2017], [K.F., B.Q. Tang, Nonlinear Analysis 2017.], [K.F. E.Latos B.Q.Tang, Annales IHP (C) 2018]

Lemma:^a For all states $\bar{\mathbf{c}} \in \mathbb{R}_{>0}^N$ satisfying $\mathcal{E}(\bar{\mathbf{c}}|\mathbf{c}_\infty) < +\infty$ and the conservation laws $\mathbb{Q}\bar{\mathbf{c}} = \mathbf{M}$, there exists a **positive constant** $H_1 = H_1(\mathbb{Q}, \mathbf{M}, \mathbf{y} \in \mathcal{C}, \mathcal{E}(\bar{\mathbf{c}}|\mathbf{c}_\infty))$ such that

$$\sum_{r=1}^{|\mathcal{R}|} \left[\sqrt{\frac{\bar{\mathbf{c}}}{\mathbf{c}_\infty}}^{\mathbf{y}_r} - \sqrt{\frac{\bar{\mathbf{c}}}{\mathbf{c}_\infty}}^{\mathbf{y}'_r} \right]^2 \geq H_1 \sum_{i=1}^N \left(\sqrt{\frac{\bar{c}_i}{c_{i,\infty}}} - 1 \right)^2.$$

Here, $\sqrt{\frac{\bar{\mathbf{c}}}{\mathbf{c}_\infty}} = \left(\sqrt{\frac{\bar{c}_1}{c_{1,\infty}}}, \dots, \sqrt{\frac{\bar{c}_N}{c_{N,\infty}}} \right)$.

This **finite-dimensional inequality** implies

$$\mathcal{D}(\mathbf{c}(t)) \geq \lambda(H_1) \mathcal{E}(\mathbf{c}(t)|\mathbf{c}_\infty),$$

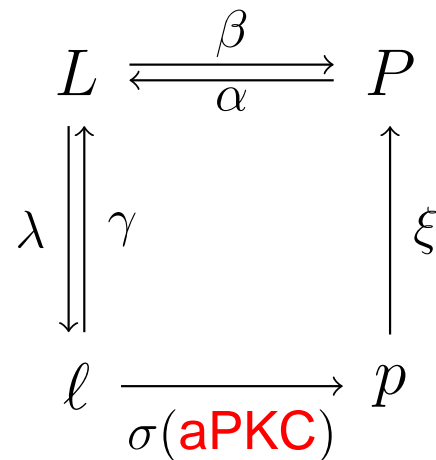
^a[K.F., B.Q. Tang, ZAMP 2018]

Systems of Reaction-Diffusion Equations

Weakly Reversible Network of Linear Reactions



Explicit exponential convergence to equilibrium state for weakly reversible volume-surface reaction-diffusion system:^a



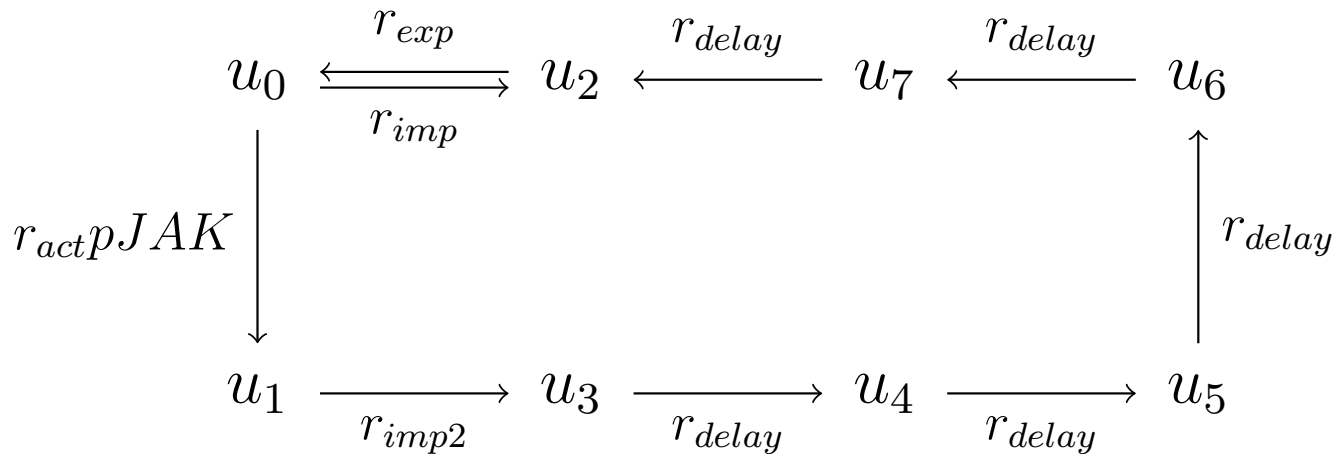
^aK.F., Bao Q. Tang Springer Proc. Math & Stats 2017

Systems of Reaction-Diffusion Equations

Weakly Reversible Network of Linear Reactions



Similar: Friedmann-Neumann-Rannacher model^a



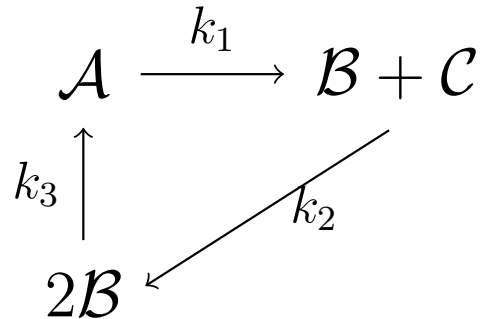
^aE. Friedmann, R. Neumann, R. Rannacher, *Well-posedness of a linear spatio-temporal model of the JAK2/STAT5 signaling pathway*, Comm. Math. Anal. 15 (2013) 76-102.

Systems of Reaction-Diffusion Equations



Boundary Equilibria

Example:



$$\begin{cases} a_t - d_a \Delta a = -k_1 a + k_3 b^2, \\ b_t - d_b \Delta b = k_1 a + k_2 bc - 2k_3 b^2, \\ c_t - d_c \Delta c = k_1 a - k_2 bc, \end{cases}$$

Boundary equilibrium $(a^*, b^*, c^*) = (0, 0, M)$.

Problem: $\mathcal{D}(a^*, b^*, c^*) = 0$, but $\mathcal{E}(c^* | c_\infty) > 0$

No global entropy-entropy dissipation estimate possible!

Boundary Equilibria

Our approach: weaker entropy-entropy dissipation estimate along solution trajectories

$$\mathcal{D}(\mathbf{c}(t)) \geq \lambda(t) \mathcal{E}(\mathbf{c}(t)|\mathbf{c}_\infty)$$

Difficulty: $\lambda(t) \rightarrow 0$ near boundary equilibria.

However, if $\lambda(t)$ satisfies $\int_0^{+\infty} \lambda(s) ds = +\infty$, then

$$\mathcal{E}(\mathbf{c}(t)|\mathbf{c}_\infty) \leq \mathcal{E}(\mathbf{c}_0|\mathbf{c}_\infty) e^{-\int_0^t \lambda(s) ds} \rightarrow 0 \quad \text{as } t \rightarrow \infty.$$

\Rightarrow (algebraic) instability of boundary equilibria

\Rightarrow Exponential convergence to positive equilibrium

Boundary Equilibria

Corresponding RD system

$$\begin{cases} a_t - d_a \Delta a = -k_1 a + k_3 b^2, & x \in \Omega, \quad t > 0, \\ b_t - d_b \Delta b = k_1 a + k_2 bc - 2k_3 b^2, & x \in \Omega, \quad t > 0, \\ c_t - d_c \Delta c = k_1 a - k_2 bc, & x \in \Omega, \quad t > 0, \\ \nabla a \cdot \nu = \nabla b \cdot \nu = \nabla c \cdot \nu = 0, & x \in \partial\Omega, \quad t > 0, \end{cases}$$

$$\inf_{x \in \Omega} b(x, t) \geq \frac{1}{\left\| \frac{1}{b_0} \right\|_{L^\infty} + 2k_3 t}, \quad \text{for all } t \geq 0.$$

Solutions would need infinite initial entropy to remain close to boundary equilibria for an unbounded time interval. ^a

^a[L. Desvillettes, K.F., B.Q. Tang, SIMA 2017]

Boundary Equilibria

Theorem:^a Let $c(t)$ be a renormalised solution of an arbitrary complex balanced network. **Assume that there exists**

$H_1 : [0, \infty) \rightarrow [0, \infty)$ such that $\int_0^\infty H_1(s) ds = +\infty$ and for a.a. $t \geq 0$

$$\sum_{r=1}^{|\mathcal{R}|} \left[\sqrt{\frac{\bar{c}(t)^{y_r}}{c_\infty}} - \sqrt{\frac{\bar{c}(t)^{y'_r}}{c_\infty}} \right]^2 \geq H_1(t) \sum_{i=1}^N \left(\sqrt{\frac{\bar{c}_i(t)}{c_{i,\infty}}} - 1 \right)^2.$$

Then, the renormalised solution $c(t)$ **converges exponentially** to the positive equilibrium c_∞ .

^a[K.F., B.Q. Tang, ZAMP 2018], [M. Pierre, T. Suzuki, Umakoshi]

Boundary Equilibria

Global Attractor Conjecture:

For any complex balanced mass action law reaction networks, all solution trajectory subject to positive initial data are conjectured to converge to the positive equilibrium c_∞ .

Proof for ODE systems by Gheorghe Craciun in 2015?

Above finite-dimensional inequality has ODE structure!?

But ODE system and averaged PDE concentrations:

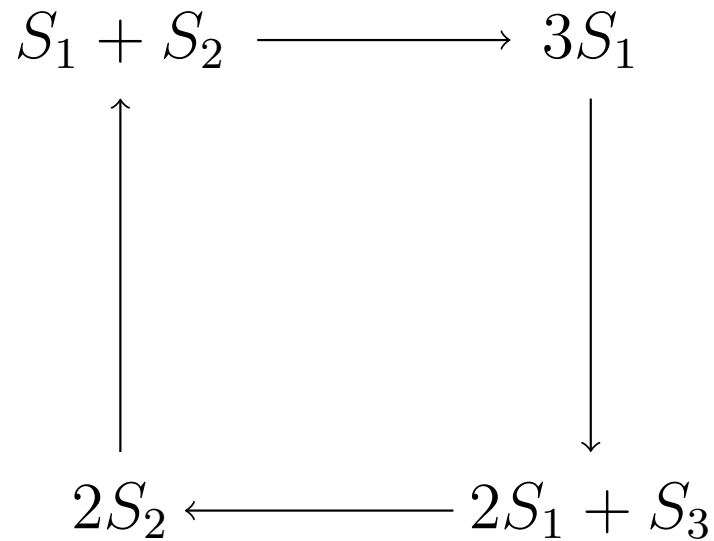
$$\frac{d}{dt} \mathbf{u} = \mathbf{R}(\mathbf{u}) \neq \overline{\mathbf{R}(\mathbf{c})} = \frac{d}{dt} \bar{\mathbf{c}}(t)$$

Systems of Reaction-Diffusion Equations



Boundary Equilibria

Boundary equilibria for complex balanced reaction networks:



Open problem!

Nonlinear Diffusion

$$\begin{cases} \partial_t c_i - d_i \Delta(c_i^{m_i}) = f_i(\mathbf{c}), & x \in \Omega, \quad t > 0, \quad i = 1, \dots, N, \\ d_i \nabla(c_i^{m_i}) \cdot \vec{n} = 0, & x \in \partial\Omega, \quad t > 0, \quad i = 1, \dots, N, \\ c_i(x, 0) = c_{i,0}(x), & x \in \Omega, \quad i = 1, \dots, N, \end{cases}$$

(i) $|f_i(\mathbf{c})| \leq C(1 + |\mathbf{c}|^\nu)$, $\forall \mathbf{c} = (c_1, \dots, c_N) \in \mathbb{R}^N$, $\forall i = 1, \dots, N$

(ii) Mass dissipation: There exist positive constants

$$\lambda_1, \dots, \lambda_N > 0 \text{ such that: } \sum_{i=1}^S \lambda_i f_i(u) \leq 0, \quad \forall \mathbf{c} \in \mathbb{R}^S$$

(iii) Quasi-positivity \Rightarrow Propagation of non-negativity

Nonlinear Diffusion

- Assume $m_i > \max\{\nu - 1; 1\}$ and $m_i > \nu - \frac{4}{d+2}$ if $d \geq 3$.

⇒ Existence of global weak nonnegative solutions

$$c_i \in C([0, \infty); L^1(\Omega)), c_i^{m_i} \in L^1(0, T; W^{1,1}(\Omega)),$$

$$f_i(\mathbf{c}) \in L^1(\Omega \times [0, T]) \text{ and}$$

$$\|c_i\|_{L^\infty(Q_T)} \leq C_T \quad \text{for all } T > 0 \quad \text{and } i = 1, \dots, N,$$

- Single reaction $\alpha_1 \mathcal{A}_1 + \dots + \alpha_M \mathcal{A}_M \xrightleftharpoons[k_f]{k_b} \beta_1 \mathcal{B}_1 + \dots + \beta_N \mathcal{B}_N$.

⇒ Exponential convergence to equilibrium $\forall 1 \leq p < \infty$,

$$\sum_{i=1}^M \|a_i(t) - a_{i\infty}\|_{L^p(\Omega)} + \sum_{j=1}^N \|b_j(t) - b_{j\infty}\|_{L^p(\Omega)} \leq C e^{-\lambda_p t}$$

Nonlinear Diffusion

Proof of existence theory extends [LP17]

- Duality estimates
- Specific bootstrap

Nonlinear Diffusion

A generalised version of Logarithmic Sobolev Inequality:

$$\int_{\Omega} \frac{|\nabla a_i|^2}{a_i^{2-m_i}} dx \geq C(\Omega, m_i) \bar{a}_i^{m_i-1} \int_{\Omega} a_i \log \frac{a_i}{\bar{a}_i} dx.$$

Degeneracy for $\bar{a}_i \sim 0$ is control by functional inequalities for **indirect diffusion effect** and **conservation law**, since not all $\bar{a}_i \sim 0$ can be small at the same time.

Setting of “slowly growing” a priori estimates:

First algebraic convergence, then exponential convergence!

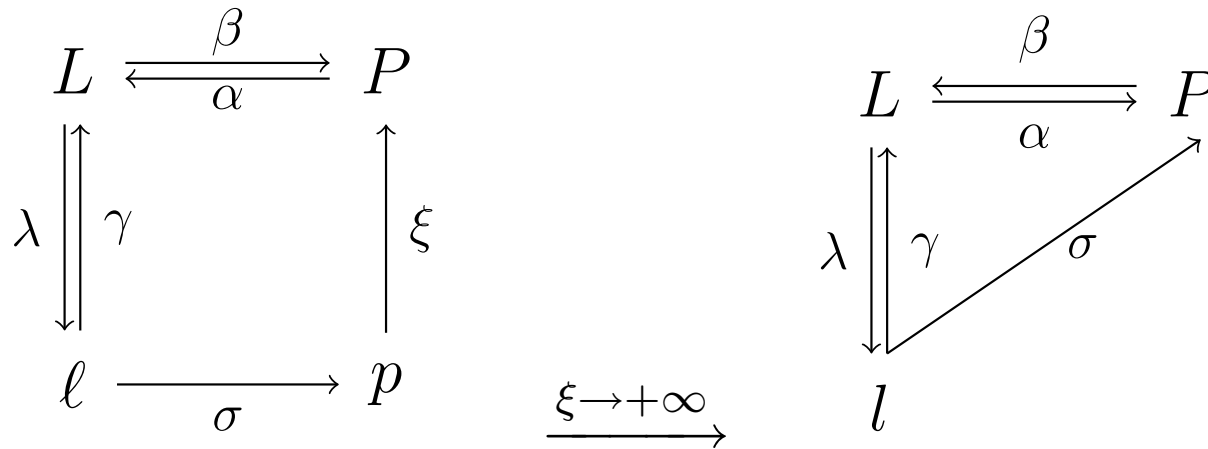
Indirect diffusion effect \sim “coercive hypocoercivity”

Quasi-steady-state approximation



QSSA as $\xi \rightarrow +\infty$

Limit of fast cortical release of Lgl: $\xi \rightarrow +\infty$



Goal: convergence towards **reduced QSSA model**.

The reduced model is **still a complex balance system**.

Quasi-steady-state approximation



QSSA as $\xi \rightarrow +\infty$

Theorem:^a

For any $(L_0, P_0, l_0, p_0) \in L^2(\Omega) \times L^2(\Omega) \times L^2(\Gamma) \times L^2(\Gamma_2)$

$$L^\xi \xrightarrow{\xi \rightarrow +\infty} L \text{ in } L^2([0, T] \times L^2(\Omega)),$$

$$P^\xi \xrightarrow{\xi \rightarrow +\infty} P \text{ in } L^1([0, T] \times \Omega),$$

$$l^\xi \xrightarrow{\xi \rightarrow +\infty} l \text{ in } L^2([0, T] \times \Gamma)$$

$$p^\xi \xrightarrow{\xi \rightarrow +\infty} 0 \text{ in } L^2([0, T] \times \Gamma_2)$$

for any $T > 0$ and up to a subsequence.

Proof: duality method [M. Pierre, D. Bothe] and entropy.

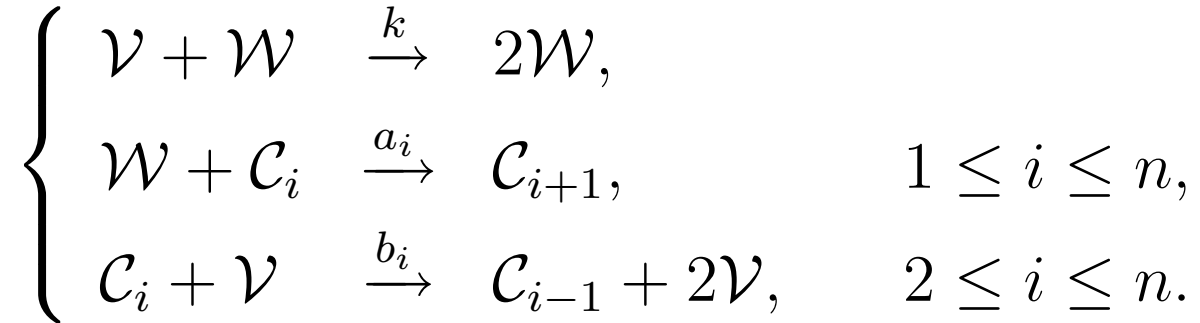
^a[T.Q.Bao, K.F., S. Rosenberger Commun. Math Sciences 2016]

Models for amyloids and protein aggregation

with Marie Doumic, Mathieu Mézache, Human Rezaei



Model for **transient oscillations** in coagulation-fragmentation experiments of PrP fibrils



Simplest two-polymer model with normalised coefficients

$$\begin{cases} \frac{dv}{dt} = v [-kw + c_2], \\ \frac{dw}{dt} = w [kv - c_1], \end{cases} \quad \begin{cases} \frac{dc_1}{dt} = -wc_1 + vc_2, \\ \frac{dc_2}{dt} = wc_1 - vc_2, \end{cases}$$

Models for amyloids and protein aggregation

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small parameter $\varepsilon = \frac{1}{k}$

$$\begin{cases} \frac{dv}{dt} = v [w_\infty - w] + \varepsilon v [v_\infty + w_\infty - v - w], \\ \frac{dw}{dt} = w [v - v_\infty] + \varepsilon w [v_\infty + w_\infty - v - w]. \end{cases}$$

Zero-order Hamiltonian $H = v_0 - v_\infty \ln v_0 + w_0 - w_\infty \ln w_0$

Full model entropy

$$\frac{d}{dt} H(v(t), w(t)) = -\varepsilon [(v - v_\infty) + (w - w_\infty)]^2.$$

\Rightarrow **Equilibration and transient oscillations** for k large.

Models for amyloids and protein aggregation

with Marie Doumic, Mathieu Mézache, Human Rezaei



THANK YOU VERY MUCH!!

